

## Another way to find the Euler-Lagrange equations system in Classical Mechanics: The Q-Gradient.

(A proposal based entirely on the Newton's 2nd law) Let be  $\vec{P} = \sum_k \vec{p}_k$ ,

$$\frac{d\vec{P}}{dt} = \vec{F}_{net}$$

$$\frac{d\vec{P}}{dt} = \vec{F}_{external} + \vec{F}_{PairBy3rdLaw}$$

$$\frac{d\vec{P}}{dt} = \vec{F}_{external} + \cancel{\vec{F}_{PairBy3rdLaw}}$$

$$\frac{d\vec{P}}{dt} - \vec{F}_{ext} = \vec{0}$$

$$\frac{d}{dt}(\vec{P}) - \vec{F}_{ext} = \vec{0}$$

$$\left\{ \frac{d}{dt}(\vec{P}) - \vec{F}_{ext} \right\} \cdot \vec{\nabla}_Q \sum \vec{r} = \vec{0} \cdot \vec{\nabla}_Q \sum \vec{r}$$

$$\frac{d}{dt}(\vec{P}) \cdot \vec{\nabla}_Q \sum \vec{r} - (\vec{F}_{ext} \cdot \vec{\nabla}_Q \sum \vec{r}) = \vec{0}$$

Since this is not the MORE-GENERAL-POSSIBLE situation about such time derivative, we guess that the following is the right choice:

$$\frac{d}{dt}(\vec{P} \cdot \vec{\nabla}_Q \sum \vec{r}) - (\vec{F}_{ext} \cdot \vec{\nabla}_Q \sum \vec{r}) = \vec{0}$$

$$\frac{d\vec{P}}{dt} \cdot \vec{\nabla}_Q \sum \vec{r} + \vec{P} \cdot \vec{\nabla}_Q \sum \frac{d\vec{r}}{dt} - (\vec{F}_{ext} \cdot \vec{\nabla}_Q \sum \vec{r}) = \vec{0}$$

Now, by a careful application of this vector equation, we should be able to find the Euler-Lagrange system of equations for any physical system already tractable in Classical Mechanics using the Lagrangian formalism. In fact, at this point we have proved the following identity: ( $\gamma$  is a constant)

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{Q}_j} \right) = \gamma \frac{d}{dt} \left( \vec{P} \cdot \frac{\partial}{\partial Q_j} \sum \vec{r} \right)$$

But wait. What is meant by  $\vec{\nabla}_Q$  as operator?... If the physical system which we are treating have  $S$  degrees of freedom, then  $\vec{\nabla}_Q$  is a first-order operator defined in a cartesian space ( $S$ -dimensional) as:

$$\vec{\nabla}_Q = \sum_{i=1}^S \frac{\partial}{\partial Q_i} \vec{e}_i$$

**Example. Dissipative forces: Air resistance**

$$\begin{aligned}
 \dot{\vec{r}} &= \begin{pmatrix} \dot{x} \\ \dot{y} \end{pmatrix} \\
 |\dot{\vec{r}}| &= \sqrt{\begin{pmatrix} \dot{x} \\ \dot{y} \end{pmatrix} \cdot \begin{pmatrix} \dot{x} \\ \dot{y} \end{pmatrix}} = \sqrt{\dot{x}^2 + \dot{y}^2} \\
 |\vec{F}_{air}| &= \beta |\dot{\vec{r}}|^n \\
 \vec{F}_{air} \cdot \dot{\vec{r}} &= |\vec{F}_{air}| |\dot{\vec{r}}| \cos(\pi) = -\beta |\dot{\vec{r}}|^{(n+1)} \\
 \frac{(\vec{F}_{air} \cdot \dot{\vec{r}})}{|\dot{\vec{r}}|} \frac{\dot{\vec{r}}}{|\dot{\vec{r}}|} &= -\beta |\dot{\vec{r}}|^{(n+1)} \frac{\dot{\vec{r}}}{|\dot{\vec{r}}|^2} = -\beta |\dot{\vec{r}}|^{(n-1)} \dot{\vec{r}}
 \end{aligned}$$

**Proportionality with the n-th power of the velocity:** ( $n > 0$ )

$$\vec{F}_{air} = -\beta (\dot{x}^2 + \dot{y}^2)^{\frac{(n-1)}{2}} \begin{pmatrix} \dot{x} \\ \dot{y} \end{pmatrix}$$

Case  $n = 1$ . Linear resistance:

$$\begin{aligned}
 \vec{F}_{air} &= -\beta \begin{pmatrix} \dot{x} \\ \dot{y} \end{pmatrix} \\
 \sum \vec{F}_{External} &= \vec{F}_{weight} + \vec{F}_{air} = -mg \begin{pmatrix} 0 \\ 1 \end{pmatrix} - \beta \begin{pmatrix} \dot{x} \\ \dot{y} \end{pmatrix} \\
 \sum \dot{\vec{r}} &= \begin{pmatrix} x \\ y \end{pmatrix} \\
 \partial_x \Sigma \dot{\vec{r}} &= \begin{pmatrix} 1 \\ 0 \end{pmatrix} \\
 \partial_y \Sigma \dot{\vec{r}} &= \begin{pmatrix} 0 \\ 1 \end{pmatrix} \\
 \partial_x \Sigma \dot{\vec{r}} &= \partial_y \Sigma \dot{\vec{r}} = \begin{pmatrix} 0 \\ 0 \end{pmatrix} \\
 \vec{P} &= m \begin{pmatrix} \dot{x} \\ \dot{y} \end{pmatrix} \\
 \vec{P} \cdot \vec{\nabla} \Sigma \dot{\vec{r}} + \vec{P} \cdot \vec{\nabla} \Sigma \dot{\vec{r}} - \vec{F}_{Ext.} \cdot \vec{\nabla} \Sigma \dot{\vec{r}} &= \vec{0} \\
 \begin{pmatrix} \ddot{x} + \frac{\beta}{m} \dot{x} \\ \ddot{y} + g + \frac{\beta}{m} \dot{y} \end{pmatrix} &= \begin{pmatrix} 0 \\ 0 \end{pmatrix}
 \end{aligned}$$